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# Cardiac re-entry dynamics and self-termination in DT-MRI based model of Human Foetal Heart

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## 2 ABSTRACT

The effect of human foetal heart geometry and anisotropy on anatomy induced drift and self-termination of cardiac re-entry is studied here in MRI based 2D slice and 3D whole heart computer simulations. Isotropic and anisotropic models of 20 weeks of gestational age human foetal heart obtained from  $100\mu\text{m}$  voxel diffusion tensor MRI data sets were used in the computer simulations. The fiber orientation angles of the heart were obtained from the orientation of the DT-MRI primary eigenvectors. In a spatially homogeneous electrophysiological monodomain model with the DT-MRI based heart geometries, cardiac re-entry was initiated at a prescribed location in a 2D slice, and in the 3D whole heart anatomy models. Excitation was described by simplified FitzHugh-Nagumo kinetics. In a slice of the heart, with propagation velocity twice as fast along the fibres than across the fibers, DT-MRI based fiber anisotropy changes the re-entry dynamics from pinned to an anatomical re-entry. In the 3D whole heart models, the fiber anisotropy changes cardiac re-entry dynamics from a persistent re-entry to the re-entry self-termination. The self-termination time depends on the re-entry's initial position. In all the simulations with the DT-MRI based cardiac geometry, the anisotropy of the myocardial tissue shortens the time to re-entry self-termination several folds. The numerical simulations depend on the validity of the DT-MRI data set used. The ventricular wall showed the characteristic transmural rotation of the helix angle of the developed mammalian heart, while the fiber orientation in the atria was irregular.

**Keywords:** cardiac arrhythmias, anatomically realistic modeling, anisotropy, anatomy induced drift, FitzHugh-Nagumo model

## 1 INTRODUCTION

Since the over a century ago hypothesis that cardiac re-entry underlies cardiac arrhythmias (Mines, 1913; Garey, 1914), and the much later confirmation of the hypothesis in cardiac tissue experiment (Allessie et al., 1973; Pertsov et al., 1993), the re-entry (*aka* spiral wave in 2D, cardiac excitation vortex in 3D), its origin and its role in sustained arrhythmias and fibrillation, as well as a possibility of its effective

control and defibrillation, have been an object of extensive theoretical study and modelling (Wiener and Rosenblueth, 1946; Balakhovsky, 1965; Krinsky, 1968; Panfilov et al., 1984; Davydov et al., 1988; Keener, 1988; Ermakova et al., 1989; Biktashev and Holden, 1994; Biktashev, 1998; Fenton and Karma, 1998; Pertsov et al., 2000; Wellner et al., 2002; Biktasheva and Biktashev, 2003; Biktashev et al., 2010, 2011; Biktasheva et al., 2015). From experiment, it is an established point of view that cardiac arrhythmias are due to a complex combination of electrophysiological (Bosch and Nattel, 2002; Workman et al., 2008; Kushiyaama et al., 2016), structural (Pellman et al., 2010; Eckstein et al., 2011; Takemoto et al., 2012; Eckstein et al., 2013), and anatomical (MacEdo et al., 2010; Anselmino et al., 2011) factors which sustain cardiac re-entry (Gray et al., 1996; Wu et al., 1998; Nattel, 2002; Yamazaki et al., 2012).

The specific effect of the heart anisotropy and anatomy on cardiac re-entry dynamics is well appreciated (Bishop et al., 2010, 2011; Bishop and Plank, 2012; Fukumoto et al., 2016), and has been studied in simplified mathematical and computer models (Fenton and Karma, 1998; Pertsov et al., 2000; Wellner et al., 2002; Rodriguez et al., 2006; Dierckx et al., 2013). The anisotropic discontinuities in the heart muscle have been commonly seen as a substrate for rise of cardiac re-entry due to the abrupt change in conduction velocity and wavefront curvature (Fenton and Karma, 1998; Spach, 2001; Smail et al., 2004). On the other hand, extensive mapping of cardiac myocyte orientation in mammalian hearts has shown that the transmural fiber arrangement, including the range of transmural change in fiber angle in ventricular wall, was consistent within a species, and varied between species (Hunter et al., 1997, p. 173). So that the anisotropic discontinuities observed in healthy hearts, and shown to be consistent within a species, might have been suspected to facilitate initiation of arrhythmias.

With the recent advance in DT-MRI technology and in High Performance Computing (HPC), the DT-MRI data sets, including anatomy and myofiber structure, can now be directly incorporated into anatomically realistic computer simulations (Vigmond et al., 2009; Bishop et al., 2010; Antonioletti et al., 2017), so that the anisotropy of the tissue in the *in-silico* model can be switched on and off for the comparison between the isotropic and anisotropic conduction, in order to see the specific anatomy effects, as well as the interplay between the anatomy and anisotropy of an individual heart.

In this paper, we present DT-MRI based anatomically and myofiber structure realistic computer simulation study of cardiac re-entry dynamics in the *in-silico* model of human foetal heart. The raw DT-MRI data (Pervolaraki et al., 2013) was segmented into the tissue/non-tissue pixels based on the MRI luminosity threshold, followed by the calculation of the fiber angles at each voxel from the diffusion-weighted DT-MRI images. The DT-MRI yields 3 eigenvalues, the primary (largest) eigenvalue was used to define local fibre orientations in the simulation study.

The main objectives are to clarify: i) whether the anatomical settings of the foetal heart might support a positive filament tension re-entry, and ii) what would it be the role of the heart anisotropy in that case. Here we demonstrate that anisotropy of the foetal heart rather facilitates self-termination of cardiac re-entry. In a 2D slice of the heart, the fiber anisotropy might change dynamics of the re-entry from pinned to anatomical re-entry. In the full 3D DT-MRI based model, depending on the location of re-entry initiation, the isotropic geometry of the heart might sustain perpetual re-entry even with a positive filament tension kinetics. While the same positive filament tension re-entry initiated at the same location of the foetal heart with the fiber anisotropy self-terminates within a fraction of the rotation period. Time of re-entry self-termination depends on the re-entry initial position. Anisotropy of the real heart speeds up re-entry self-termination. The geometry and anisotropy of the heart together ensure the fastest self-termination of cardiac re-entry.

The novel significance of our findings is that we demonstrate that the heart anisotropy might have rather anti-arrhythmic function as it facilitates fast self-termination of cardiac re-entry. A general role of fiber anisotropy in the genesis and sustenance of arrhythmias could be addressed by numerics even on idealised and simplified geometries with different spatial distributions of anisotropy. The biomedical question addressed in the manuscript is whether self-terminating ventricular arrhythmias can occur in a developing foetal heart, as has inferred from fECG data in (Benson et al., 2015).

## 2 MATERIALS AND METHODS

### 2.1 DT-MRI data sets

Tissue acquisition followed medical termination of pregnancy with written and informed consent, and Ethical approval from Lothian Research Ethics Committee (reference 08/S1101/1). Temporary storage of the tissue for imaging was in premises licensed under the UK 2004 Human Tissues Act.

The DT-MRI data set used in this study was of a 143 days gestational age (DGA) human foetal heart described in (Pervolaraki et al., 2013). It was selected as by 143 days the smooth, transmural  $120^\circ$  transmural rotation in helix fiber angle is well established (Pervolaraki et al., 2013; Mekkaoui et al., 2013; Pervolaraki et al., 2017). The heart was immersed in formalin shortly after dissection, and imaged in formalin after two weeks in formalin.

MRI acquisition was performed in a Bruker Biospin (Ettlingen, Germany) 9.4 Tesla vertical NMR/S System with a 22mm imaging coil for Hydrogen (1 H). A three-dimensional diffusion weighted spin echo sequence was carried out at  $20^\circ\text{C}$  with  $0.1\text{mm}^2$  resolution, echo time = 15ms, repetition time = 500ms, with 6 averages and a b value of  $1000\text{s/mm}^2$ . In each scan, diffusion weighted images were obtained in 12 directions. The average scan time was 24 hours.

### 2.2 DT-MRI based anatomy model

Figure 1 shows a cross section of the 143 days of gestational age (DGA) foetal heart, with the already formed intramural myolaminar structure, and yet a bit irregular surface epicardial, endocardial, and septal fibers, see also Figure 4 in (Pervolaraki et al., 2013, p. 5) for the color-encoded fractional anisotropy (FA) and all the three components of the fiber angles in the human foetal heart. While in an adult heart, pinning of cardiac re-entry to endocardium structures such as pectinate muscles junction with crista terminalis had been previously reported (Wu et al., 1998; Yamazaki et al., 2012; Kharche et al., 2015a). The DT-MRI based foetal heart model offered a unique opportunity to see whether the 20 weeks of gestation age intramural heart structure was capable to support cardiac re-entry, because at that foetal development stage it would not be possible yet for the re-entry to pin to the endocardium fine features, for these anatomical structures were yet to be developed later.

The DT-MRI data sets of the  $128 \times 128 \times 128$  voxels size, with voxel resolution of  $\sim 100\mu\text{m}$ , of 143 days of gestational age (DGA) human foetal heart (Pervolaraki et al., 2013), were converted into the BeatBox (Antonioletti et al., 2017) regular Cartesian mesh .bbg geometry format, containing the DT-MRI Cartesian coordinates of the heart tissue points together with the corresponding components of the diffusion tensor primary eigenvectors (Antonioletti et al., 2017). The .bbg file is an ASCII text file, each line in which describes a point in a regular mesh in the following format:

$x, y, z, \text{status}, \text{fibre}_x, \text{fibre}_y, \text{fibre}_z$





**Figure 1. The 143 DGA human foetal heart (Pervolaraki et al., 2013).** BeatBox (Antonioletti et al., 2017) geometry .bbg format visualisation with projections of the unit vectors of the local fibre orientation onto the cross-section plane: a smooth intramural fibre distribution has developed, with some surface irregular fibers from earlier developmental stages, as seen in Figure 4 (Pervolaraki et al., 2013, p. 5).

105 Here  $x$ ,  $y$ ,  $z$  are integer Cartesian coordinates of a DT-MRI voxel, `status` is a flag with a nonzero-  
 106 value for a tissue point, and `fibre_x`, `fibre_y`, `fibre_z` are  $x$ -,  $y$ - and  $z$ -components of the fibre  
 107 orientation vector at that point. To reduce the size of the .bbg files, only the tissue points, that is points  
 108 with nonzero `status` need to be specified, because the BeatBox solver will ignore the void points with  
 109 zero `status` in any case. Although the original DT-MRI images data sets had  $128 \times 128 \times 128$  voxels size,  
 110 the actual dimensions of the foetal heart minimum bounding box were  $67 \times 91 \times 128$ , with 181070 tissue  
 111 points.

112 The raw DT-MRI anatomy data (Pervolaraki et al., 2013) were segmented into the “tissue”/“not tissue”  
 113 pixels discretion based on the MRI luminosity threshold, with the Cartesian fiber angles at each voxel  
 114 obtained from the diffusion-weighted DT-MRI images. Only this basic segmentaion of the raw DT-MRI  
 115 anatomy data (Pervolaraki et al., 2013) was taken into account in the computer simulation of cardiac  
 116 re-entry dynamics, so we shall refer to it as the raw DT-MRI based anatomy model.

117 Any raw MRI and micro-CT image data show tiny bits of the preparation tissue, which usually get into  
 118 the image together with the heart at the preparation stage. In the case of the raw DT-MRI image of the foetal  
 119 heart (Pervolaraki et al., 2013), there happened to be a tiny bit of tissue at the bottom of the MRI image,  
 120 adjacent to the apex of the foetal heart, see Figure 1, and the original Figure 2(a), last panel, in (Pervolaraki  
 121 et al., 2013, p. 3). In order to see whether this tiny “leftover” piece of the heart tissue is capable to affect  
 122 the outcome of a re-entry simulation, we edited the MRI images in order to remove the “leftover” piece,  
 123 and then considered a comparison of the re-entry simulation in the two 3D DT-MRI based heart models:  
 124 i) raw DT-MRI based anatomy model, and ii) “edited” DT-MRI based anatomy model without the tiny  
 125 “leftover” piece of the heart tissue.

126 In case of the 2D model of a slice of the heart, in order to construct the 2D diffusivity tensor, the fibres  
 127 vectors were projected into the plane of the slice of the heart.

## 128 2.3 Cardiac tissue model

To investigate the effects of anatomy on cardiac re-entry dynamics we used *monodomain* tissue model with non-flux boundary conditions

$$\begin{aligned}\frac{\partial \mathbf{u}}{\partial t} &= \mathbf{f}(\mathbf{u}) + \nabla \cdot \hat{\mathbf{D}} \nabla \mathbf{u}, \\ \vec{n} \cdot \hat{\mathbf{D}} \nabla \mathbf{u} \Big|_G &= 0,\end{aligned}\tag{1}$$

where  $\mathbf{u}(\vec{r}, t) = (u, v)^T$ ,  $\vec{r}$  is the position vector,  $\mathbf{f}(\vec{r}, t) = (f, g)^T$  is the FitzHugh-Nagumo (Winfree, 1991) kinetics column-vector

$$\begin{aligned}f(u, v) &= \alpha^{-1}(u - u^3/3 - v), \\ g(u, v) &= \alpha(u + \beta - \gamma v),\end{aligned}\tag{2}$$

129 with the parameter values  $\alpha = 0.3$ ,  $\beta = 0.71$ ,  $\gamma = 0.5$ , which in an infinite excitable medium support a  
130 rigidly rotating vortex with positive filament tension (Biktashev et al., 1994). The simplified FHN model  
131 was intentionally chosen for this study in order to fully eliminate the possible effects of a realistic cell  
132 excitation kinetics, such as *e.g.* meander (Winfree, 1991), alternans (Karma, 1994), negative filament  
133 tension (Biktashev et al., 1994), etc., and in order to enhance and highlight the pure effects of the heart  
134 anatomy and anisotropy on the cardiac re-entry outcome.  $\hat{\mathbf{D}} = \mathbf{Q}\hat{\mathbf{P}}$ , where  $\mathbf{Q} = \text{diag}(1, 0) = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$  is  
135 the matrix of the relative diffusion coefficients for  $u$  and  $v$  components, and  $\hat{\mathbf{P}} = [P_{jk}] \in \mathbb{R}^{3 \times 3}$  is the  $u$   
136 component diffusion tensor, which has only two different eigenvalues: the bigger, simple eigenvalue  $P_{\parallel}$   
137 corresponding to the direction along the tissue fibers, and the smaller, double eigenvalue  $P_{\perp}$ , corresponding  
138 to the directions across the fibres, so that

$$P_{jk} = P_{\perp} \delta_{jk} + (P_{\parallel} - P_{\perp}) f_j f_k,\tag{3}$$

139 where  $\vec{f} = (f_k)$  is the unit vector of the fiber direction;  $\vec{n}$  is the vector normal to the tissue boundary  $G$ .  
140 In the isotropic simulation,  $P_{\parallel}$  and  $P_{\perp}$  values were fixed at  $P_{\parallel} = P_{\perp} = 1$  (corresponding 1D conduction  
141 velocity 1.89, in the dimensionless units of Eqs. (1)-(2)). In the anisotropic simulations,  $P_{\parallel}$  and  $P_{\perp}$  values  
142 were fixed at  $P_{\parallel} = 2$ ,  $P_{\perp} = 0.5$  (corresponding conduction velocities 2.68 and 1.34 respectively). All the  
143 conduction velocities have been computed for the period waves with the frequency of the free spiral wave  
144 in the model, i.e. 11.36. With the isotropic diffusivity ( $P_{\parallel} = P_{\perp} = 1$ ) equal to the geometric mean between  
145 the faster and the slower anisotropic diffusivities ( $P_{\parallel} = 2$ ,  $P_{\perp} = 1/2$ ), the isotropic conduction velocity  
146 1.89 was almost exactly the same as the geometric mean  $\approx 1.895$  of the faster and slower (2.68 and 1.34  
147 respectively) anisotropic conduction velocities, chosen in order to minimize the maximal relative difference  
148 between the isotropic and anisotropic propagation speeds.

149 All the computer simulations presented here were done using the BeatBox (Antonioletti et al., 2017)  
150 software package with the explicit time-step Euler scheme, on the Cartesian regular grid with space step  
151 discretization  $\Delta x = 0.1$ , time step discretisation  $\Delta t = 0.001$ ; 5-point stencil for isotropic, and 9-point  
152 stencil for anisotropic Laplacian approximation in 2D simulations; 7-point stencil for isotropic, and 27-  
153 point stencil for anisotropic Laplacian approximation in 3D simulations. The re-entry was initiated by the  
154 phase distribution method (Biktashev and Holden, 1998): in the 2D simulations, at a prescribed location of

the cross section of the DT-MRI based anatomical model shown in Figure 1; in the 3D simulations, at a prescribed location of the full DT-MRI based whole heart anatomical model.

The FHN model Eqs. (1)-(2) is not biophysically detailed and is formulated in dimensionless units. So, for illustrative purposes only, if we took the FHN time unit to be 1 t.u. = 40 ms this would give the APD90 of 125.6 ms which is within the range reported e.g. by (Zhu et al., 2016). The spiral wave period is then 454 ms: at our chosen kinetics parameters, the spirals in FHN model have a big core. With the space step discretization  $\Delta x = 0.1$  in FHN simulations, whereas the real grid resolution is  $100 \mu\text{m} = 0.1 \text{ mm}$ , hence we would have the FHN space unit 1 s.u. = 1 mm. The (geometric mean) diffusivity of 1 s.u.<sup>2</sup>/t.u. therefore works out as  $0.025 \text{ mm}^2/\text{ms}$ , and the corresponding conduction velocity of 1.89 s.u./t.u. is  $0.04725 \text{ mm/ms}$ . That gives a spiral wavelength of  $\approx 21.45 \text{ mm}$ . The conduction velocity thus obtained is within the range reported in (Pervolaraki et al., 2013), whereas the diffusivity is about four times smaller than the one used in (Pervolaraki et al., 2014). We must stress here that, since the FHN kinetics is not biophysically detailed, one should not expect anything more than an order-of-magnitude correspondence with reality.

### 3 RESULTS

#### 3.1 2D MRI-based “slice” simulations

In the 2D simulations, Figure 2, a counter-clockwise re-entry was initiated by the phase distribution method (Biktashev and Holden, 1998), with the initial center of rotation placed at the prescribed location  $x_0 = 40, y_0 = 60$  in the 2D cross section of the DT-MRI based anatomical model shown in Figure 1.

In the Figure 2(a-b), it can be seen that in both isotropic and anisotropic 2D simulations, at  $t = 0$ , there was identical location of the initial re-entry rotation center: roughly in the middle of the slice, in the vicinity of the septum cuneiform opening.

Figure 2(a) shows *isotropic* dynamics of the re-entry, that is with the fiber orientation data “turned OFF”, so that only the geometry of the isotropic homogeneous slice affects the re-entry. While in an infinite medium the chosen FHN parameter values  $\alpha = 0.3, \beta = 0.71, \gamma = 0.5$  produce rigidly rotating spiral (Winfree, 1991), the boundaries of the foetal heart slice model cause the drift of the re-entry. The re-entry does not terminate because of the reflection from the inexcitable boundaries (Biktashev and Holden, 1994), but after the transient first rotation around the septum cuneiform opening, the tip of the re-entry firmly pins (at  $t = 35$ ) to the sharp lower end of the cuneiform opening, see Figure 2(a).

Figure 2(b) shows *anisotropic* dynamics of the re-entry, that is with the fiber orientation data “turned ON”, so that both the geometry and the anisotropy of the otherwise homogeneous slice of the heart affect the dynamics of the re-entry, causing its drift. In the anisotropic slice, the re-entry also does not terminate at the inexcitable boundaries, but after a faster than in the previous isotropic case transient, compare the time units labels in the Figs. 2(a-b), the anisotropy of the medium turns the initial spiral wave into the fast anatomical re-entry around the septum cuneiform opening, see Figure 2(b).

#### 3.2 3D whole heart MRI-based simulations

##### 3.2.1 Raw DT-MRI anatomy model

In the 3D whole heart *raw* DT-MRI based simulations shown in the Figure 3 and Figure 4, a counter-clockwise excitation vortex was initiated by the phase distribution method (Biktashev and Holden, 1998), with the initial position of the transmural vortex filament (yellow line) at the prescribed location *along the*

194  $x$  axis at  $y_0 = 40, z_0 = 60$ . It can be seen in Figure 3 isotropic, and Figure 4 anisotropic 3D simulations  
 195 that, at  $t = 0$ , there was identical initial location of the filament of the excitation vortex: that is transmurally,  
 196 roughly in the middle through the ventricles of the heart.

197 Figure 3 shows the *isotropic dynamics* of the excitation vortex, that is with the fiber orientation data  
 198 “turned OFF”, so that only the geometry of the otherwise isotropic homogeneous foetal heart affects  
 199 dynamics of the vortex. At the chosen parameter values  $\alpha = 0.3, \beta = 0.71, \gamma = 0.5$ , the FHN vortex has  
 200 positive filament tension (Biktashev et al., 1994), and, depending on topology of an isotropic homogeneous  
 201 medium, either collapses or straightens up between parallel boundaries. In the isotropic simulations of  
 202 the foetal heart, boundaries of the heart cause the vortex to drift and collapse. However, there exist initial  
 203 locations of the excitation vortex, which although result in the drift of the vortex, still do not lead to the  
 204 expected collapse of the vortex with positive filament tension. One of such outcomes is shown in the  
 205 Figure 3. Here, following the geometry of the heart, after a very short transient, the initial vortex filament  
 206 breaks into the two major pieces, each of which finds its own synchronous pathway in the “isotropic”  
 207 foetal heart, resulting in the seemingly perpetual cardiac re-entry, which failed to self-terminate within the  
 208 extended simulation time  $t = 30$ . The time course of the instant number of filaments (blue dashed line) is  
 209 shown in Figure 12(a), and the time course of the instant total length of the filaments (blue dashed line)  
 210 is shown in Figure 12(b). It can be seen that, after a very fast transient increase in both the number of  
 211 filaments and the total length of the filaments, these fail to disappear, and keep oscillating around above zero  
 212 constants. The maximum instant number of the filaments in the simulation was 9 at  $t = 1.0$ , the maximum  
 213 instant total length of the filaments was 127.1 at  $t = 1.6$ ; while the time average number of the filaments in  
 214 the simulation was 3.4, and the time average total length of the filaments was 29.3, see the summary of the  
 215 simulation detail in Figure 11.

216 Figure 4 shows *anisotropic dynamics* of the excitation vortex, that is with the fiber orientation data “turned  
 217 ON”, so that both the geometry and the anisotropy of the otherwise homogeneous model of the foetal  
 218 heart affect dynamics of the vortex. Here, the anisotropy of the heart causes fast transient distortion of the  
 219 filament, and drift towards the inexcitable boundary of the heart, followed by the very fast self-termination  
 220 of the vortex by  $t = 5.3$ . The time course of the instant number of filaments (pink dotted line) is shown in  
 221 Figure 12(a), and the time course of the instant total length of the filaments (pink dotted line) is shown in  
 222 Figure 12(b). It can be seen that, after a very fast transient increase in both the number of the filaments  
 223 and the total length of the filaments, all the filaments rapidly disappear. The maximum instant number of  
 224 the filaments was 13 at  $t = 0.8$ , that is higher and achieved faster than in the isotropic conduction. The  
 225 maximum instant total length of the filaments was 179.7 at  $t = 1.8$ , again much higher than in the isotropic  
 226 conduction. The time average number of the filaments in the simulation was 6.5, twice higher than in the  
 227 isotropic conduction, and the time average total length of the filaments was 91.2, three times higher than in  
 228 the isotropic conduction, see the summary of the simulation detail in Figure 11.

229 In the 3D whole heart *raw* DT-MRI based simulations shown in the Figure 5 and Figure 6, a counter-  
 230 clockwise excitation vortex was initiated by the phase distribution method (Biktashev and Holden, 1998),  
 231 with the initial position of the transmural vortex filament (yellow line) at the prescribed location *along the*  
 232  $y$  axis at  $x_0 = 40, z_0 = 60$ , that is *perpendicular* to the initial orientation of the vortex filament shown in  
 233 Figure 3 and Figure 4. It can be seen in Figure 5 isotropic, and in Figure 6 anisotropic 3D simulations, that  
 234 at  $t = 0$ , there was identical initial location of the filament of the excitation vortex: that is transmurally,  
 235 roughly in the middle through the ventricles of the foetal heart, and *perpendicular* to the initial orientation  
 236 of the vortex filament shown in Figure 3 and Figure 4.

Figure 5 shows the *isotropic dynamics* of the excitation vortex, that is with the fiber orientation data “turned OFF”, so that only the geometry of the otherwise isotropic homogeneous foetal heart affects dynamics of the vortex. Here, contrary to the expectation for the positive filament tension vortex to always contract, the organising filament first transiently extends intramurally along the tissue walls, before the final break up into the two ring-like pieces, each of which then quickly contracts and terminates at the opposite base and apex regions of the heart by  $t = 4$ . The time course of the instant number of filaments (blue dashed line) is shown in Figure 12(c), and the time course of the instant total length of the filaments (blue dashed line) is shown in Figure 12(d). It can be seen that, after a very fast transient increase in both the number of the filaments and the total length of the filaments, all the filaments rapidly disappear. The maximum instant number of the filaments was 12 at  $t = 1.4$ , with the maximum instant total length of the filaments 188.0 achieved at  $t = 1.5$ . The time average number of the filaments in the simulation was 4.7, and the time average total length of the filaments was 80.5, see the summary of the simulation detail in Figure 11. It can be seen that in this isotropic conduction simulation, with just a different location/orientation of the initial vortex filament, the time average total length of the vortex filaments was two and a half times higher than in the failed to self-terminate isotropic conduction simulation shown in Figure 3.

Figure 6 shows *anisotropic dynamics* of the excitation vortex, that is with the fiber orientation data “turned ON”, so that both the geometry and the anisotropy of the otherwise homogeneous model of the heart affect dynamics of the vortex leading to its really fast termination at the apex of the heart by  $t = 2.6$ . The time course of the instant number of filaments (pink dotted line) is shown in Figure 12(c), and the time course of the instant total length of the filaments (pink dotted line) is shown in Figure 12(d). It can be seen that, after a very fast transient increase in both the number of the filaments and the total length of the filaments, all the filaments rapidly disappear. The maximum instant number of the filaments was 17 at  $t = 0.5$ , that is higher and achieved faster than in the corresponding isotropic conduction shown in Figure 5. The maximum instant total length of the filaments was 278.6 at  $t = 0.9$ , again much higher than in the corresponding isotropic conduction. The time average number of the filaments in the simulation was 9.5, twice higher than in the isotropic conduction, and the time average total length of the filaments was 152.6, also twice higher than in the isotropic conduction, see the summary of the simulation detail in Figure 11.

It can be seen in the *raw* DT-MRI model simulations shown in Figure 3, Figure 4, Figure 5, and Figure 6, that although the organising filament of the vortex could not get through into the accidental “leftover” piece of tissue adjacent to the apical region, the piece still got activated, and could have served as an artificial “capacitor” affecting dynamics of the re-entry. In order to check whether this might be the case, we edited the original *raw* DT-MRI model by removing in the MRI the foreign “leftover” piece, and repeated the whole heart isotropic and anisotropic simulations from the same two orthogonal initial locations of the re-entry, similar to the shown in Figure 3, Figure 4, Figure 5, and Figure 6.

### 3.2.2 “Edited” DT-MRI anatomy model

In the 3D whole heart “*edited*” MRI model simulations shown in the Figure 7 and Figure 8, a counter-clockwise excitation vortex was initiated by the phase distribution method (Biktashev and Holden, 1998), with the initial position of the transmural vortex filament (yellow line) at the prescribed location *along the*  $x$  axis at  $y_0 = 40$ ,  $z_0 = 60$ . It can be seen in Figure 7 isotropic, and in Figure 8 anisotropic 3D simulations, that, at  $t = 0$ , there was identical initial location of the filament of the excitation vortex: that is transmurally, roughly in the middle through the ventricles of the foetal heart, similar to the initial location of the vortex filament in the *raw* DT-MRI simulations shown in Figure 3 and Figure 4.

Figure 7 shows *isotropic dynamics* of the vortex, that is with the fiber orientation data “turned OFF”, so that only the geometry of the otherwise isotropic homogeneous foetal heart affects dynamics of the vortex. Here, following the geometry of the heart, the initial filament also breaks into the two major pieces, each of which also finds its own synchronous pathway similar to the beginning of the raw DT-MRI simulation shown in Figure 3. However, this time, after just a few rotations, the two re-entries find their end in their almost synchronous termination in the base region of the heart by  $t = 16.9$ . The time course of the instant number of filaments (green dashed line) is shown in Figure 12(a), and the time course of the instant total length of the filaments (green dashed line) is shown in Figure 12(b). It can be seen that, after a very fast transient increase in both the number of the filaments and the total length of the filaments, all the filaments rapidly disappear. The maximum instant number of the filaments was 9 at  $t = 0.2$ . The maximum instant total length of the filaments was 122.3 at  $t = 0.7$ . The time average number of the filaments in the simulation was 3.2, and the time average total length of the filaments was 28.7, see the summary of the simulation detail in Figure 11. It can be seen that, in the “edited” MRI isotropic simulation in Figure 7, the maximum instant and the time average number of the filaments, as well as the maximum instant and the time average total length of the filaments, were practically the same as in the corresponding raw MRI isotropic simulation with the failed to self-terminate re-entry shown in Figure 3. The only quantitative difference between the two isotropic simulations, *i.e.* the perpetual re-entry in Figure 3 vs the self-termination in Figure 7, was that, in the “edited” MRI isotropic simulation, without the “leftover” piece of tissue adjacent to the apex of the heart, the maximum instant number of the filaments, and the maximum instant total length of the filaments were achieved much faster: by  $t = 0.2$  and  $t = 0.7$  correspondingly.

Figure 8 shows *anisotropic dynamics* of the vortex, that is with the fiber orientation data “turned ON”, so that both the geometry and the anisotropy of the otherwise homogeneous model of the heart affect dynamics of the vortex. Here, the anisotropy of the heart also causes significant transient distortion of the initial filament, followed by the fast drift towards the apex, with the ultimate termination at the AV border before a completion of a single rotation, very similar to the corresponding raw DT-MRI anisotropic simulation shown in Figure 4. However, in the “edited” MRI model without the “leftover” piece adjacent to the apex, repolarisation of the heart is faster than it was in the presence of the “incidental capacitor” effect in the corresponding raw DT-MRI simulation shown in Figure 4. The time course of the instant number of filaments (red solid line) is shown in Figure 12(a), and the time course of the instant total length of the filaments (red solid line) is shown in Figure 12(b). It can be seen that, after a very fast transient increase in both the number of the filaments and the total length of the filaments, all the filaments rapidly disappear by  $t = 4.8$ . The maximum instant number of the filaments was 17 at  $t = 0.4$ , that is higher than in the corresponding raw anisotropic conduction shown in Figure 4, and twice higher than in the corresponding “edited” isotropic conduction shown in Figure 7. The maximum instant total length of the filaments was 180.3 at  $t = 0.7$ , similar to the corresponding raw anisotropic conduction shown in Figure 4, and much higher than in the corresponding “edited” isotropic conduction shown in Figure 7. The time average number of the filaments in the simulation was 6.3, similar to the corresponding raw anisotropic conduction shown in Figure 4, and twice higher than in the corresponding “edited” isotropic conduction shown in Figure 7. The time average total length of the filaments was 95.4, similar to the corresponding raw anisotropic conduction shown in Figure 4, and three times higher than in the corresponding raw isotropic conduction shown in Figure 4. Time to the maximum number of the filaments ( $t = 0.4$ ), and time to the maximum total length of the filaments ( $t = 0.7$ ), were similar to the corresponding “edited” isotropic conduction shown in Figure 7, and twice faster than in the corresponding raw anisotropic conduction shown in Figure 4, see the summary of the simulation detail in Figure 11. So, that the anisotropy of the heart increased the number and the total



length of the filaments, and shortened the time to self-termination by folds, while the “incidental capacitor” effect of the “leftover” piece adjacent to the apex slowed down the process.

In the 3D whole heart “*edited*” MRI simulations shown in Figure 9 and Figure 10, a counter-clockwise excitation vortex was initiated by the phase distribution method (Biktashev and Holden, 1998), with the initial position of the transmural filament (yellow line) at the prescribed location *along the y axis* at  $x_0 = 40, z_0 = 60$ . It can be seen in Figure 9 isotropic, and in Figure 10 anisotropic 3D simulations, that at  $t = 0$ , there was the identical initial location of the filament: that is transmurally, roughly in the middle through the ventricles of the heart, perpendicular to the filament initial location in the “*edited*” MRI simulations shown in Figure 7 and Figure 8, and similar to the initial location of the filament in the raw DT-MRI simulations shown in Figure 5 and Figure 6.

Figure 9 shows “*edited*” MRI *isotropic dynamics* of the vortex, that is with the fiber orientation data “turned OFF”, so that only the geometry of the otherwise isotropic homogeneous foetal heart affects the vortex. Here, again contrary to the expectation for a positive filament tension vortex to always contract, the organising filament first transiently extends intramurally before breaking up into the two ring-like pieces, each of which quickly contracts and terminates at the opposite base and apex regions of the heart, identical to what was seen in the raw DT-MRI simulation shown in Figure 5. The time course of the instant number of filaments (green dashed line) is shown in Figure 12(c), and the time course of the instant total length of the filaments (green dashed line) is shown in Figure 12(d). It can be seen that, after a very fast transient increase both in the number of the filaments and in the total length of the filaments, all the filaments rapidly disappear by  $t = 4.0$ . The maximum instant number of the filaments was 12 at  $t = 0.5$ . The maximum instant total length of the filaments was 190.7 at  $t = 1.5$ . The time average number of the filaments in the simulation was 4.7, and the time average total length of the filaments was 82.9, see the summary of the simulation detail in Figure 11. So, from the comparison with the corresponding raw MRI isotropic simulation shown in Figure 5, it seems that, for that particular initial location of the filament, the “leftover” tissue did not play any role in the re-entry self-termination time.

Figure 10 shows *anisotropic dynamics* of the vortex, that is with the fiber orientation data “turned ON”, so that both the geometry and the anisotropy of the otherwise homogeneous model of the heart affect the vortex, which, in the absence of the “incidental capacitor” effect of the “leftover” piece, results in the fastest possible termination of the re-entry at the apex of the heart by  $t = 2.3$ , before the vortex first rotation ever started. The re-entry termination time here is more than twice shorter than in both the raw MRI isotropic simulation shown in Figure 5, and in the “*edited*” isotropic simulations shown in Figure 9; shorter than in the corresponding simulation with the “incidental capacitor” effect shown in the Figure 6, and times shorter than in any of the simulations of the re-entry with the perpendicular initial location of the filament shown in the Figure 3, Figure 4, Figure 7, and Figure 8. So, that the main reasons for the re-entry fastest self-termination seem to be the initial location of the filament and the anisotropy of the heart.

In Figure 11, we have summarized the results of the *raw* MRI simulations shown in Figure 3, Figure 4, Figure 5, Figure 6, and the “*edited*” MRI simulations shown in Figure 7, Figure 8, Figure 9, and Figure 10. The re-entry termination time, in the time units of Eqs. (1)-(2), is shown under each respected whole heart model initiation cite panel. It can be seen that the anisotropy of the heart causes at least twice faster termination of re-entry. It also can be seen that indeed the “leftover” piece of tissue connected to the apical region of the heart in the raw DT-MRI model served as an artificial “capacitor” affecting the re-entry dynamics, and significantly prolonged the life time of re-entry initiated at the particular location/orientation respective to the “capacitor”.

366 Finally, the 3D anatomically realistic simulations of the foetal heart show that the anisotropy of the heart  
367 causes the fast transient increase in the number and the total length of the filaments, see Figure 12, with the  
368 typical fast drift towards the apex area of the heart, and re-entry self-termination, see also the movies in the  
369 Supplementary Material section 5.1.

## 4 DISCUSSION

370 Although the role of heart anatomy and anisotropy in the origin and sustainability of cardiac arrhythmias  
371 has been appreciated for a long time, there is limited experimental evidence to clarify detail of the heart  
372 anatomy effect on persistent cardiac arrhythmias and fibrillation. In particular, the theoretically plausible  
373 hypothesis that the anisotropic discontinuities in the heart might be a source of rise for cardiac re-entry due  
374 to the abrupt change in conduction velocity and wavefront curvature (Fenton and Karma, 1998; Spach,  
375 2001; Smaill et al., 2004) was in controversy with the observation that the transmural fiber arrangement,  
376 including the range of transmural change in fiber angle in ventricular wall, although varied between  
377 species (Hunter et al., 1997, p. 173), was consistent within a species. So that the question was that, if the  
378 pro-arrhythmic mechanism of cardiac re-entry initiation by the anisotropic discontinuities in a heart (Fenton  
379 and Karma, 1998; Spach, 2001; Smaill et al., 2004) was correct, what would then have been a reason for the  
380 consistent structure (Hunter et al., 1997, p. 173) of the anisotropic discontinuities in healthy mammalian  
381 hearts. The combination of the High Performance Computing with the high-resolution DT-MRI based  
382 anatomy models of the heart allows anatomically realistic *in-silico* testing of the effects of individual  
383 heart anatomy and anisotropy on the cardiac re-entry dynamics (Vigmond et al., 2009; Bishop et al., 2010;  
384 Kharche et al., 2015a,b; Antonioletti et al., 2017). In this paper, for the first time, we present the anatomy  
385 and myofiber structure realistic computer simulation study of the cardiac re-entry dynamics in the DT-MRI  
386 based model of the human foetal heart (Pervolaraki et al., 2013).

387 The human foetal heart single anisotropic geometry dataset used in these simulations needs to be a  
388 typical and accurate representation of the cardiac structure and tissue architecture, with a spatial resolution  
389 suitable for numerical solution of the excitation equations. For a smoothly changing anisotropic geometry  
390 the imaged dataset may be interpolated to provide sufficient spatial resolution. DT-MRI has provided  
391 17 – 200 cubic voxel datasets of human foetal heart (Pervolaraki et al., 2012, 2013, 2017; Tang et al.,  
392 2016). Finite difference numerical solutions of propagation in biophysically detailed models of adult  
393 cardiac tissue need  $\sim 100\mu m$  space steps in a Cartesian coordinate system. During foetal development  
394 the structure of the heart changes, the ventricular wall is compactified and trabeculated, and the size of  
395 the heart increases. Any atlas of the developing human foetal heart would need to come from longitudinal  
396 studies (which are impractical and potentially unethical), or from a large number of hearts obtained and  
397 imaged at different gestational ages. Here, all simulations are done on a single cardiac geometry that was  
398 selected as the transmural changes in fiber helix angle had developed. The results illustrated are specific  
399 for this particular anisotropic cardiac geometry, which is critically determined by how it was imaged and  
400 reconstructed. The fiber orientation angles of the heart were obtained from the orientation of the DT-MRI  
401 primary eigenvectors. Propagation velocity was twice as fast along the fibres than across the fibers. In all  
402 the simulations on this DT-MRI based cardiac geometry, the anisotropy of the myocardial tissue shortens  
403 the duration of re-entry by several fold. The numerical simulations depend on the validity of the DT-MRI  
404 dataset used. The ventricular wall showed the characteristic transmural rotation of the helix angle of the  
405 developed mammalian heart, while the fiber orientation in the atria was irregular. We expect the results be  
406 subsequently verified on other anatomy data, including different technique data e.g. MRI vs DT-MRI vs  
407 micro-CT vs serial section histology, etc.

408 The comparative isotropic vs anisotropic simulation of the otherwise homogeneous model of the foetal  
409 heart shows that, in the 2D slice of the heart, the fiber anisotropy might change the re-entry dynamics from  
410 the re-entry pinned at the sharp end of the septum cuneiform opening, Figure 2(a), into the fast anatomical  
411 re-entry around the opening, Figure 2(b). Note that, despite of the only basic segmentation of the MRI  
412 model into the tissue/not tissue points, and the ventricles not being isolated from the atria, because of the  
413 2D re-entry pinning to either the sharp end of the septum opening, as in the isotropic simulation shown  
414 in Figure 2(a), or to the whole septum opening, as in the case of anatomical re-entry in the anisotropic  
415 simulation shown in Figure 2(b), the tip of the re-entry never penetrated from the ventricles into the atria.

416 Although, from the cardiac physiology point of view, the only basic segmentation of the raw DT-MRI  
417 data (Pervolaraki et al., 2013) into the tissue/non-tissue pixels might be seen as a major limitation of the  
418 study, from the non-linear science point of view, the use of the raw MRI data, as an example of a nature  
419 provided medium to study a re-entry dynamics, gives an important insight into the pure anatomy induced  
420 drift in an otherwise homogeneous 2D medium, and into the possibility of pinning of the re-entry not to a  
421 major blood vessel but to a sharp end of an anatomical opening (Biktasheva et al., 2015); and into that a  
422 real fiber anisotropy is capable to turn the pinned re-entry into an anatomical one. Importantly though, the  
423 2D simulations in Figure 2 are an important step to highlight the role and the necessity of the whole heart  
424 structure in the re-entry dynamics and self-termination.

425 In the 3D DT-MRI based *isotropic* model of the foetal heart, depending on the initial location/orientation  
426 of the re-entry organising filament, the geometry of the foetal heart might sustain perpetual cardiac re-entry  
427 even with a positive filament tension, see Figure 3 and Figure 12(a-b). However, if the same positive  
428 filament tension vortex is initiated at the exactly same location/orientation in the *anisotropic* DT-MRI based  
429 model, the fiber structure of the foetal heart facilitates fast self-termination of cardiac re-entry, see Figure 4  
430 and Figure 12(a-b).

431 From the respective comparison of the “isotropic vs anisotropic” simulations in Figure 3 vs Figure 4,  
432 and Figure 7 vs Figure 8, it can be seen that, whereas the re-entry organizing filaments were capable to  
433 penetrate from the ventricles to atria in the isotropic simulations shown in Figure 3 and Figure 7, the abrupt  
434 change in the fiber angles between the atria and the ventricles, which one can see in Figure 1, did not allow  
435 the organizing filaments to penetrate from ventricles into atria in the anisotropic simulations shown in  
436 Figure 4 and Figure 8, so that the anisotropy of ventricles could complete the speedy elimination of the  
437 re-entry within a single rotation.

438 The comparison of re-entry termination times in the whole heart raw DT-MRI simulations shown in  
439 Figure 3, Figure 4, Figure 5, and Figure 6, with the corresponding “edited” MRI simulations shown  
440 in Figure 7, Figure 8, Figure 9, and Figure 10, showed that, although the filament of cardiac re-entry  
441 never got through into the tiny piece of the “leftover” tissue adjacent to the apex of the heart in the raw  
442 DT-MRI, the “not heart” excitable tissue served as an accidental “capacitor”, significantly prolongating  
443 the life time of cardiac re-entry initiated at a particular location/orientation respective to the “leftover”  
444 piece location/orientation. The comparison of the re-entry termination times in Figure 11 and Figure 12  
445 shows that the higher maximum number of the filaments and the bigger total length of the filaments tend to  
446 correlate with the faster termination of re-entry. However, the usual transient increase in both the number  
447 of the filaments and in the total length of the filaments, still failed to identify and terminate the persistent  
448 re-entry shown in Figure 3 and Figure 12(a-b).

449 The “isotropic vs anisotropic” comparison of self-termination time, both in the raw and in the “edited”  
450 MRI whole heart simulations, confirmed that, regardless of with or without the “leftover” piece adjacent

to the apex, anisotropy of the heart shortens re-entry self-termination time several folds, see Figure 11 and Figure 12. Figure 12 shows that anisotropy increases the maximum number and the maximum total length of the filaments. The bigger maximum number and the maximum total length of the filaments tend to correlate with the faster termination of re-entry. The biggest transient total lengths of the filaments was in the anisotropic case of re-entry initiated *along the y axis*, see the corresponding dotted pink and solid red lines in panel (d) in Figure 12, which ensured the re-entry fastest termination. It can be seen from Figure 5, Figure 6, Figure 9, and Figure 10, that the initial position of the filament *along the y axis* allowed the filament to grow intramurally, thus maximally increasing the transient total length of the filaments, and shortening their termination time.

The comparison of the “edited” MRI simulations of thus completely isolated heart shown in Figure 7, Figure 8, Figure 9, and Figure 10, with the raw DT-MRI simulations shown in Figure 3, Figure 4, Figure 5, Figure 6, provide an important new insight into cardiac re-entry dynamics. Namely, that an excitable tissue accidentally adjacent to the heart might serve a capacitor capable to extend re-entry self-termination time, see for the respective comparison Figure 3 vs Figure 7, Figure 4 vs Figure 8, Figure 5 vs Figure 9, and Figure 6 vs Figure 10, summarised in Figure 11 and Figure 12. The latter suggests a possible new mechanism for a persistent cardiac re-entry. That is if, apart from the major blood vessels normally adjacent to the heart in vivo, and affecting cardiac re-entry dynamics, there were also an accidental “touching” of the heart by an adjacent excitable tissue, for example, due to a change of posture in the night (Hu et al., 2011), the “incidental capacitor” effect could have extended re-entry self-termination time, up to the failure to self-terminate. The “incidental capacitor” hypothesis could be an additional explanation to the circadian rhythm (Jeyaraj et al., 2012), for the elusive and difficult to reproduce longer arrhythmia episodes reported in the night ECGs as opposed to the on average shorter arrhythmias in the day time ECGs. Although the raw DT-MRI simulations with the “leftover” piece of tissue might have been seen a limitation of the study, the real heart in vivo does not exist in complete isolation from the main blood vessels and other neighboring tissues. So, we believe that our “incidental” leftover tissue results only once more confirm the importance and the necessity of taking into account of the real anatomical settings and surrounding of the heart for the full appreciation of cardiac re-entry dynamics.

The BeatBox DT-MRI based *in-silico* comparative study confirms the heart anatomy and anisotropy functional effect on cardiac re-entry self-termination as opposed to its sustainability, pinning to anatomical features, transformation from pinned to anatomical re-entry, while the anisotropy of the tissue facilitates re-entry self-termination.

One of the limitations of the present study is the use of the simplified FitzHugh-Nagumo (Winfree, 1991) excitation model, chosen for this study in order to eliminate the effects of realistic cell excitation kinetics, such as *e.g.* meander (Winfree, 1991), alternans (Karma, 1994), negative filament tension (Biktashev et al., 1994), etc., and enhance and highlight the pure effects of the heart anatomy and anisotropy on the re-entry. The realistic cell excitation models should be used in the future studies, in order to clarify specific interplay of the cell kinetics with the heart anatomy and anisotropy.

As it can be seen from Figure 1 (for the color-encoded fractional anisotropy (FA) and for the color-encoded all the three components of the fiber angles see also Figure 4 in (Pervolaraki et al., 2013, p. 5)), formation of the foetal heart fiber structure at the epicardium and endocardium is not completed yet, so that only the already formed intramural laminar myofibers affect dynamics of the re-entry. Although the use of the not fully formed foetal heart can be seen as a limitation of the study, on the other hand, the irregular epicardium and endocardium fibers seem to prevent a re-entry from pinning to the fine endocardium features, which were yet to be developed (Pervolaraki et al., 2013) later on. We appreciate

that the anatomy and fiber structure differences in the foetal and in a fully formed adult heart might alter the re-entry dynamics, such as in *e.g.* the reported case of re-entry pinning to the junction of pectinate muscles with crystae terminalis in adult human atrium (Wu et al., 1998; Yamazaki et al., 2012; Kharche et al., 2015a). That is, although it is possible to initiate a cardiac re-entry in the tiny 1.4g (at 143 DGA) foetal heart (Pervolaraki et al., 2013), the already formed intramural laminar anisotropy of the foetal heart facilitates the re-entry self-termination, Figure 11. With the hindsight of the present study, in an adult heart, because of the pinning opportunities provided by the endocardium anatomical features (Wu et al., 1998; Yamazaki et al., 2012; Kharche et al., 2015a), there must exist additional mechanisms to facilitate cardiac re-entry self-termination (Clayton et al., 1993).

The most serious limitation of the study is that only the basic segmentation of the raw DT-MRI data (Pervolaraki et al., 2013) into the tissue/non-tissue pixels based on the MRI luminosity threshold, and only the primary eigenvalues of fibres orientation, were taken into account in the BeatBox (Antonioletti et al., 2017) computer simulations of cardiac re-entry dynamics. Further levels of DT-MRI segmentation, in order to take into account *e.g.* the heart collagen skeleton, isolation of ventricles from atria, etc., will inevitably change the outcome of the re-entry, by adding the electrically impermeable barriers to cardiac re-entry. Currently, this further segmentation is added into DT-MRI based models via complex rule based image post-processing (Lombaert et al., 2012; Gahm et al., 2013), which brings in artificial assumptions, as well as limits the number of available segmented data sets. From the non-linear science point of view we pursued in this initial study, the rationale was to use the raw DT-MRI as an example of a nature provided medium to study a re-entry dynamics. In the future, the multichannel computer tomography might offer automatic tissue segmentation, so that multi-level segmented DT-MRI based models become more available, and be used in the BeatBox (Antonioletti et al., 2017) anatomically and biophysically realistic simulations.

Finally, we believe that a simple “mechanistic” explanation, although often craved for, might be rather inadequate/premature here, and will require better theoretical understanding of the demonstrated potential effect of the heart anatomy and anisotropy on cardiac re-entry dynamics, for it is not a particular feature, or a sequence of features, but rather the whole complex of the shape, anisotropy, and position of the heart within the body, which affects the re-entry dynamics in a particular way, and which seems to have evolved in order to ensure the fastest self-termination of cardiac re-entry. If our hypothesis is correct, it might explain the difficulties with reproducibility of the arrhythmia in vivo and in an isolated heart. A general role of fiber anisotropy in the genesis and sustenance of arrhythmias could be and has been addressed by numerics on idealised and simplified geometries with different spatial distributions of anisotropy (Fenton and Karma, 1998). The biomedical question addressed in the manuscript is whether self-terminating ventricular arrhythmias can occur in a developing foetal heart, as has inferred from fECG data in (Benson et al., 2015). The most important novel finding of the paper is that, contrary to what currently seems to be a commonly accepted view of the pro-arrhythmic nature of cardiac anisotropy, the point of view based on the mainly theoretical and simplified anatomy models studies, for the first time ever, and for the first time in a real whole heart DT-MRI based model, we have demonstrated that anisotropy of the heart might have rather anti-arrhythmic effect, as it facilitates the fastest self-termination of cardiac re-entry.

## 5 NOMENCLATURE

### 5.1 Resource Identification Initiative

BeatBox, RRID:SCR\_015780



## CONFLICT OF INTEREST STATEMENT

536 The authors declare that the research was conducted in the absence of any commercial or financial  
537 relationships that could be construed as a potential conflict of interest.

## ETHICS STATEMENT

538 Tissue acquisition followed medical termination of pregnancy with written and informed consent, and  
539 Ethical approval from Lothian Research Ethics Committee (reference 08/S1101/1). Temporary storage of  
540 the tissue for imaging was in premises licensed under the UK 2004 Human Tissues Act.

## AUTHOR CONTRIBUTIONS

541 IVB, AVH, EP and FCW contributed conception and design of the study; RAA, AVH and EP contributed  
542 the human foetal heart DT-MRI data sets; IVB and FCW converted the DT-MRI data sets into BeatBox  
543 geometry format and ran the simulations; IVB performed visualisation of the simulations and analysis, and  
544 wrote the first draft of the manuscript. FCW approved submission of the paper. IVB, RAA, AVH, and EP  
545 contributed to manuscript revision, read and approved the submitted version.

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## SUPPLEMENTARY MATERIAL

552 The Supplementary Material for this article can be found online at ...

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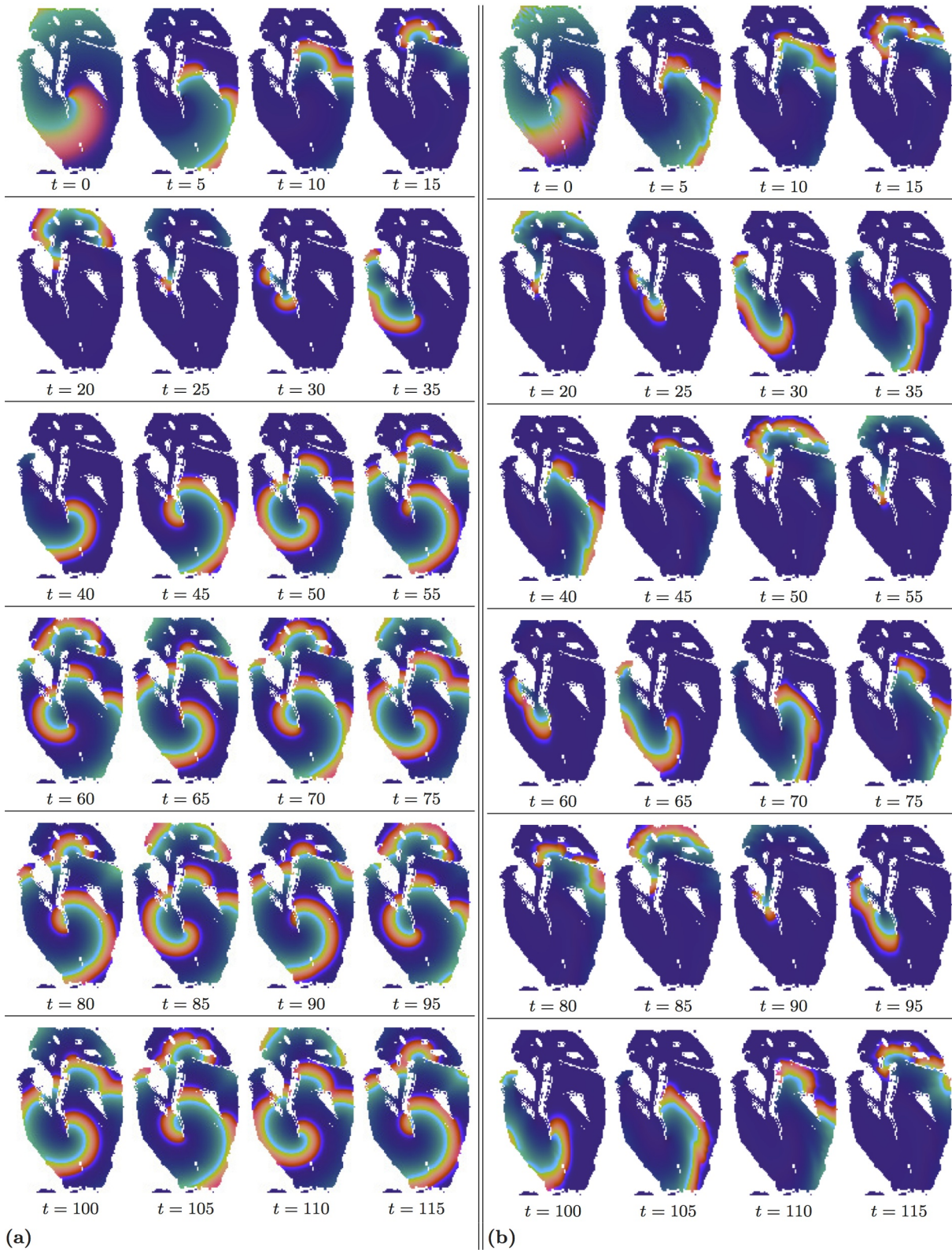
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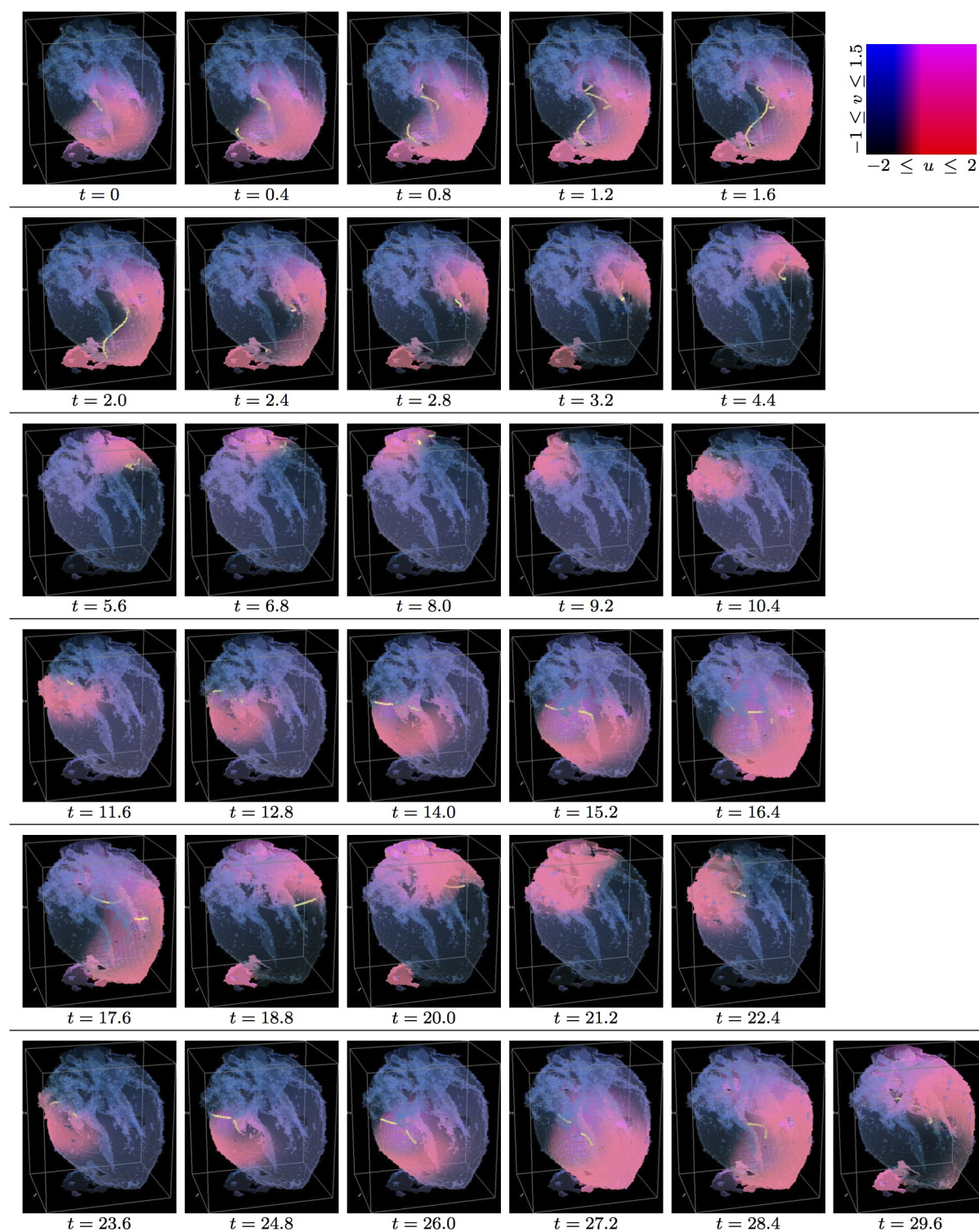
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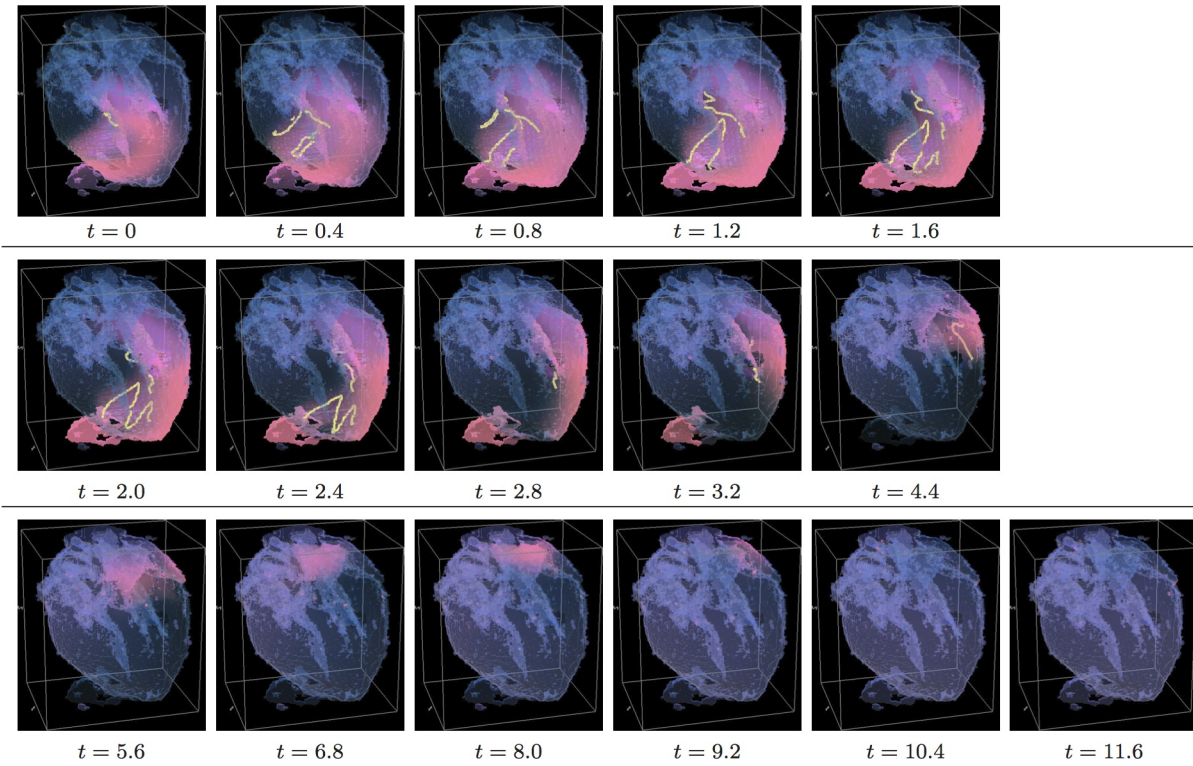


**Figure 2. Anisotropy effect in the 2D slice simulations,** time shown under each panel in time units of Eqs. (1)-(2). **a) Isotropic conduction:** after the transient first rotation around the septum cuneiform opening, the slow excitation re-entry pins to the sharp low end of the opening in the foetal heart. (Multimedia view) Fig2a.mpg. **b) Anisotropic conduction:** after the fast transient first rotation around the septum cuneiform opening, the anisotropy of the foetal heart turns the initial spiral wave into the fast anatomical re-entry around the septum cuneiform opening. (Multimedia view) Fig2b.mpg.

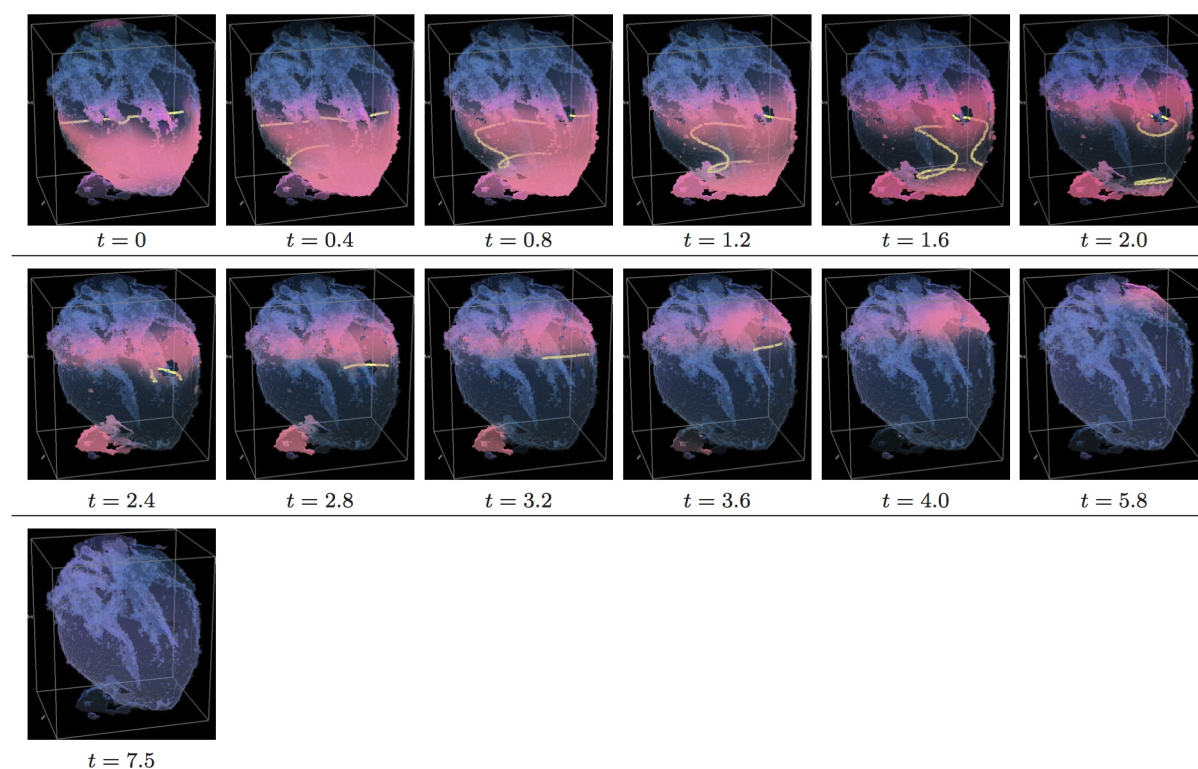


**Figure 3. Isotropic whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). After a short transient the organising filament of the initial vortex breaks into the two short pieces each of which finds its own synchronous perpetual pathway, resulting in the perpetual cardiac re-entry in the foetal heart. (Multimedia view) Fig3.mpg.

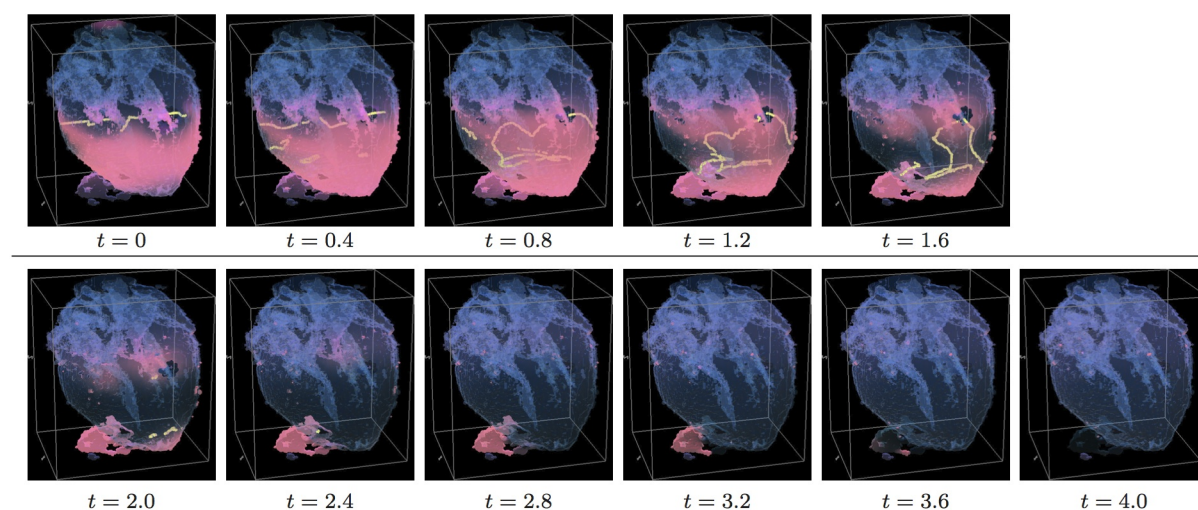




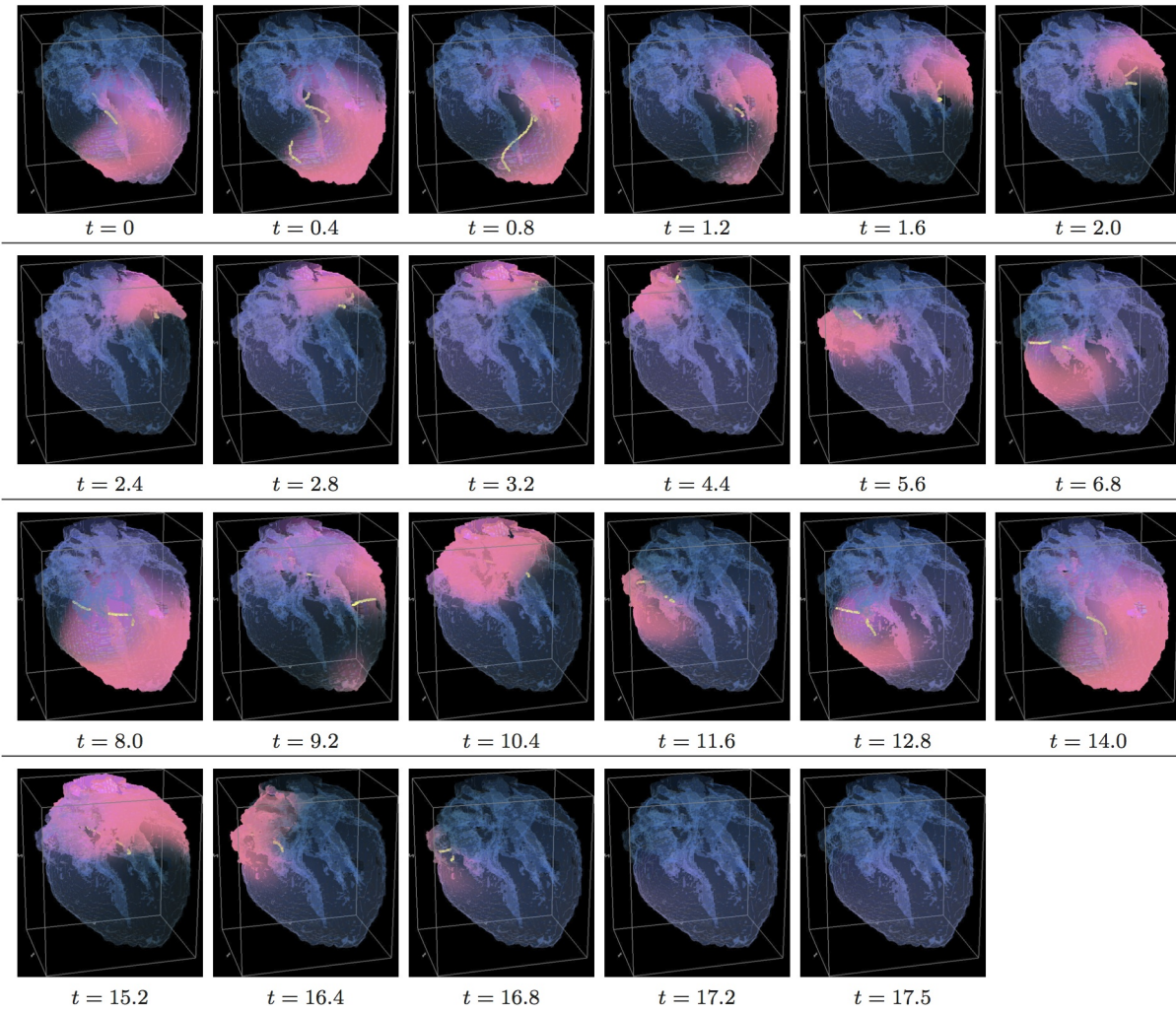
**Figure 4. Anisotropic whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). The anisotropy of the heart causes the fast transient distortion of the organising filament of the initial excitation vortex and drift towards the inexcitable boundary of the heart, ultimately resulting in the very fast self termination of the excitation vortex. (Multimedia view) Fig4.mpg.



**Figure 5. Isotropic whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). After a short transient the organising filament of the initial vortex breaks into the two pieces each of which fast terminates: one at the base and another at the apex of the heart. (Multimedia view) Fig5.mpg.

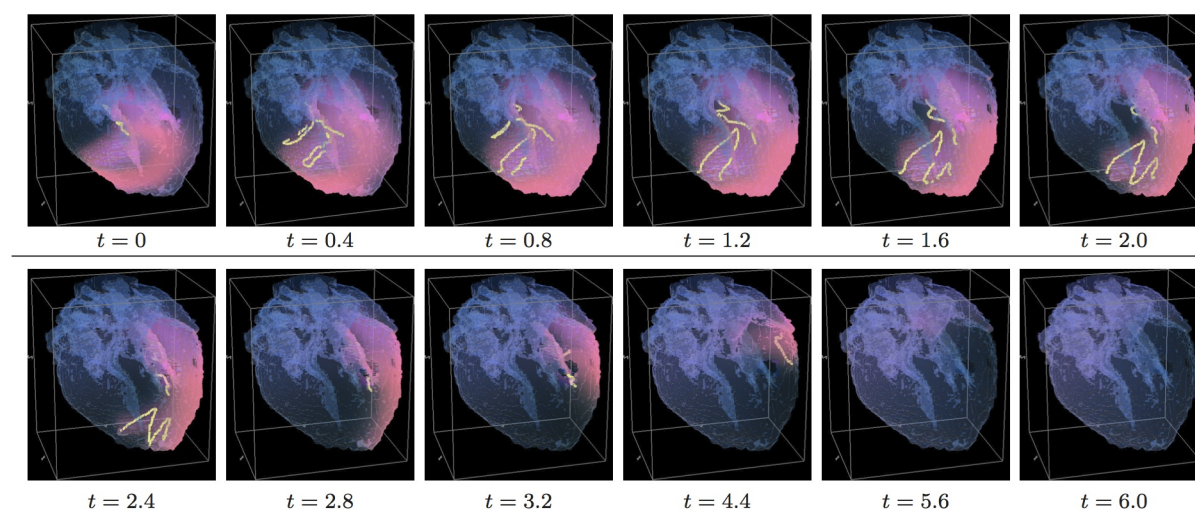


**Figure 6. Anisotropic whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). The anisotropy of the heart causes the fast transient distortion of the organising filament of the initial excitation vortex, followed by the fast drift and self-termination at the apex of the heart. (Multimedia view) Fig6.mpg.

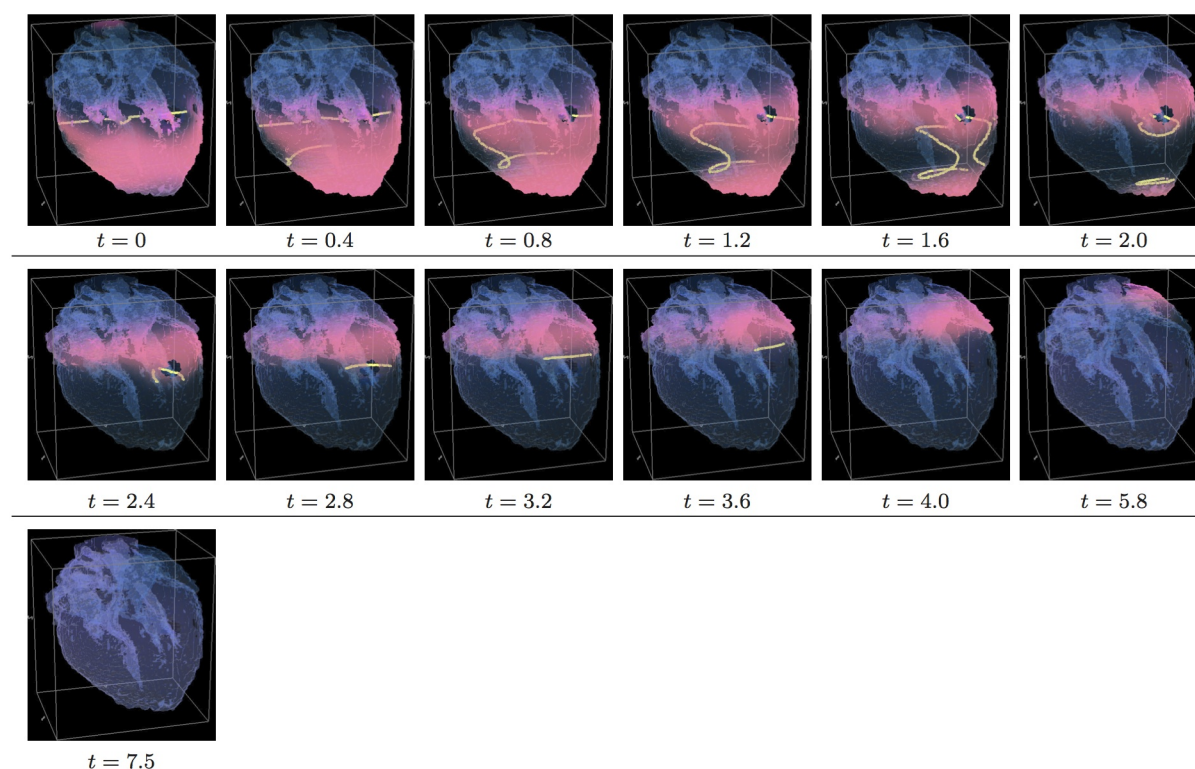


**Figure 7. Isotropic “edited” whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). After a short transient the organising filament of the initial vortex breaks into the two short pieces each of which finds its own synchronous pathway, resulting after a few rotations in the synchronous termination of the filaments in the base of the foetal heart. (Multimedia view) Fig7.mpg.

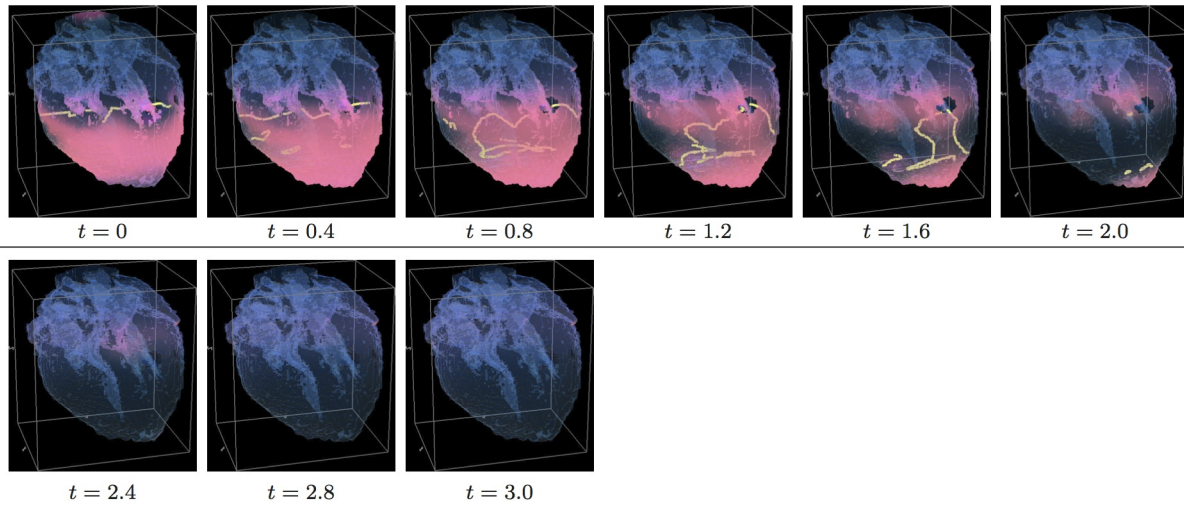




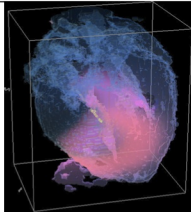
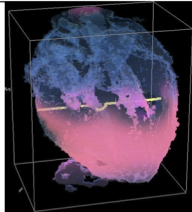
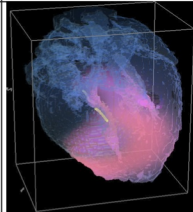
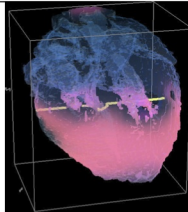
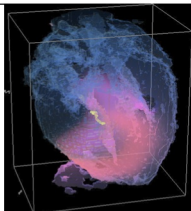
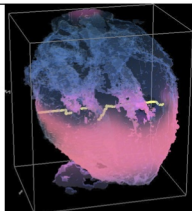
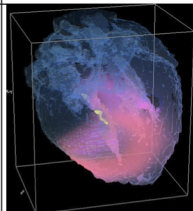
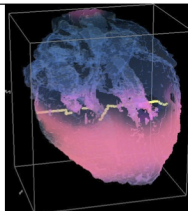
**Figure 8. Anisotropic “edited” whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). The anisotropy of the heart causes the significant transient distortion of the organising filament of the initial vortex, followed by its fast drift towards the apex and the ultimate termination before completing a single rotation. (Multimedia view) Fig8.mpg.



**Figure 9. Isotropic “edited” whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). After a short transient the organising filament of the initial vortex breaks into the two pieces each of which fast terminates: one at the base and another at the apex of the heart. (Multimedia view) Fig9.mpg.

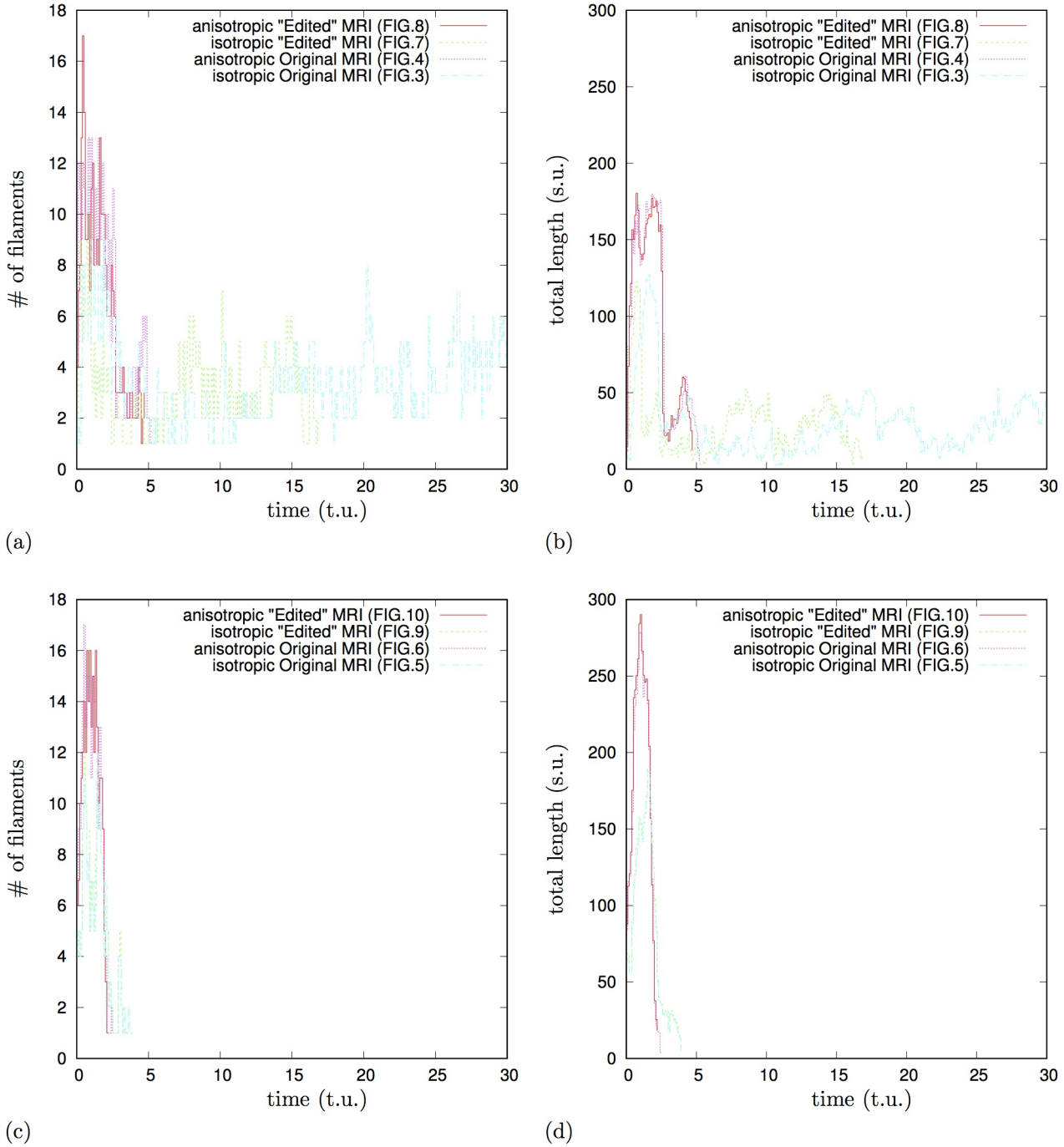


**Figure 10. Anisotropic “edited” whole heart simulation.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices; time shown under each panel in time units of Eqs. (1)-(2). The anisotropy of the heart causes the fast significant transient distortion of the organising filament of the initial excitation vortex, followed by the fast drift towards the apex and ultimate termination before the first rotation has ever started. (Multimedia view) Fig10.mpg.

	Original MRI		“Edited” MRI		
Isotropic					
	Simulation FIG.	FIG.3	FIG.5	FIG.7	FIG.9
	Filaments Termination time (t.u)	$t=\infty$	$t=4.0$	$t=16.9$	$t=4.0$
	Recovery time (t.u)	$t=\infty$	$t=7.5$	$t=17.5$	$t=7.5$
	Average # of filaments	3.4	4.7	3.2	4.7
	Average total length	29.3	80.5	28.7	82.9
	$max(\# \text{ of filaments})(t_{max(\# \text{ of filaments})})$	9 ( $t=1.0$ )	12 ( $t=1.4$ )	9 ( $t=0.2$ )	12 ( $t=0.5$ )
	$max(\text{total length})(t_{max(\text{total length})})$	127.1 ( $t=1.6$ )	188.0 ( $t=1.5$ )	122.3 ( $t=0.7$ )	190.7 ( $t=1.5$ )
AnIsotropic					
	Simulation FIG.	FIG.4	FIG.6	FIG.8	FIG.10
	Filaments Termination time (t.u)	$t=5.3$	$t=2.6$	$t=4.8$	$t=2.3$
	Recovery time (t.u)	$t=10.0$	$t=4.0$	$t=6.0$	$t=3.0$
	Average # of filaments	6.5	9.5	6.3	10.3
	Average total length	91.2	152.6	95.4	177.0
	$max(\# \text{ of filaments})(t_{max(\# \text{ of filaments})})$	13 ( $t=0.8$ )	17 ( $t=0.5$ )	17 ( $t=0.4$ )	16 ( $t=0.7$ )
	$max(\text{total length})(t_{max(\text{total length})})$	179.7 ( $t=1.8$ )	278.6 ( $t=0.9$ )	180.3 ( $t=0.7$ )	290.2 ( $t=1.0$ )

**Figure 11. Whole heart simulation: re-entry termination times.** The translucent foetal heart is shown in blue, excitation front shown in red (see the color box in Figure 3), the yellow lines are the instant organising filaments of the excitation vortices. Re-entry self-termination time in time units of Eqs. (1)-(2) is shown under each simulation Figure3–Figure10 initiation panel. Comparison of the respective isotropic (top row) vs anisotropic (bottom row) simulations shows that, regardless of with or without the “leftover” piece, anisotropy results in faster termination of re-entry, and at least twice shorter recovery time. Respective comparison of the original MRI with the corresponding “edited out leftover” simulations shows that the leftover “incidental capacitor” effect, depending of the re-entry location/orientation with respect to the “incidental capacitor” own location/orientation, might significantly prolongate cardiac re-entry life time. The bigger number and the total length of the filaments tend to correlate with the faster termination of re-entry, though these fail to identify persistent re-entry in Figure 3 simulation.





**Figure 12. Whole heart simulation: time course of the number of filaments #, and of the total length of the filaments, time and the total length of the filaments shown in the time and space units of Eqs. (1)-(2). Initial position of the vortex filament along the *x* axis, Figure 3, Figure 4, Figure 7, and Figure 8 simulations: (a) time course of the number of filaments # ; (b) time course of the total length of the filaments. Initial position of the vortex filament along the *y* axis, Figure 5, Figure 6, Figure 9, and Figure 10 simulations: (c) time course of the number of filaments # ; (d) time course of the total length of the filaments. Anisotropy increases the transient number and the transient total length of the filaments. The bigger transient number and the total length of the filaments tend to correlate with the faster termination of re-entry. The biggest transient total lengths of the filaments was in case of the re-entry initiated along the *y* axis, panel (d), which ensured its fastest termination. It can be seen from Figure 5, Figure 6, Figure 9, and Figure 10, that the initial position of the filament along the *y* axis allowed it to grow intramurally, thus maximally increasing the transient total length of the filaments, and speeding up their termination.**